

Hamiltonian Monodromy: Dynamical Systems and Physics

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Description of the system

▷ $\mathbb{R}^4 \rightsquigarrow (\mathbf{q}, \mathbf{p}) = ((x, y), (p_x, p_y))$

▷ Hamiltonian differential equation

$$H: \mathbb{R}^4 \rightarrow \mathbb{R} \text{ s.t. } \dot{q} = \frac{\partial H}{\partial p}, \dot{p} = -\frac{\partial H}{\partial q}$$

Property:

$\varphi(t)$

$$\frac{dH(\varphi(t))}{dt} = 0 \text{ (First integral)}$$

Solutions contained in the level curves of H

▷ $K: \mathbb{R}^4 \rightarrow \mathbb{R}$ first integral independent of H

- Completely integrable system

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Arnold-Liouville Theorem

The connected components of a regular compact fiber of \mathcal{EM} are diffeomorphic to real tori \mathbb{T}^2 . Moreover, the diffeomorphism is such that

$$\begin{aligned}\psi: U &\rightarrow D^2 \times \mathbb{T}^2 \\ (q, p) &\mapsto (I, \phi)\end{aligned}$$

and the differential equations transform into

$$\begin{aligned}\dot{I} &= 0 \\ \dot{\phi} &= V(I)\end{aligned}$$

A-A coordinates \Rightarrow solutions, linear flow over invariant tori.

- Is this change of coordinates global?

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The Spherical Pendulum

- Video

- ◇ Differential equations

$$q = (x, y, z) \in \mathbb{S}^2, p = (p_x, p_y, p_z) \in T_q\mathbb{S}^2$$

$$\dot{q} = \vec{p}$$

$$\dot{p} = -\vec{e}_z - (p_x^2 + p_y^2 + p_z^2 - z)\vec{q}$$

- ◇ Hamiltonian (Energy)

$$H = \frac{1}{2}(p_x^2 + p_y^2 + p_z^2) + z$$

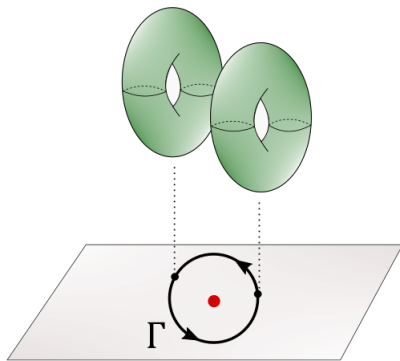
- ◇ First integral independent of H

$$K = xp_y - yp_x$$

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Hamiltonian Monodromy (HM)

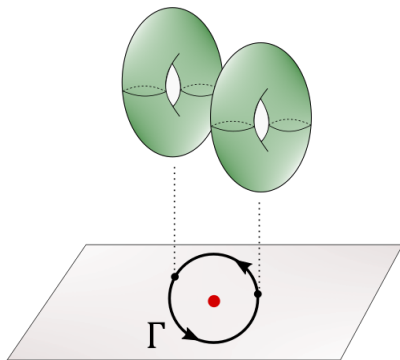
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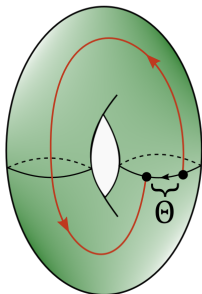


- Non trivial HM \Rightarrow non global A-A coordinates.

Rotation number Θ

Symmetry under rotations associated to $K \Rightarrow \Theta(h, k)$.

$$\varepsilon \mathcal{M}^{-1}(h, k) \subset \mathbb{R}^4$$

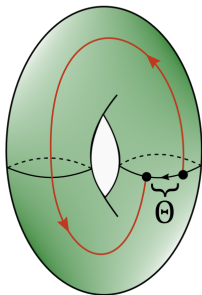


In this case, HM is completely determined by Θ : $\mathcal{M}_\Gamma = \begin{pmatrix} 1 & \frac{\Delta_\Gamma \Theta}{2\pi} \\ 0 & 1 \end{pmatrix}$

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Spectral Lax Pairs

\mathbb{R}^4

- ▷ Lax pair: L, M , 2×2 complex matrices time dependent s.t.

$$\dot{L} = [M, L] = ML - LM$$

- ▷ Spectral Lax pair: $L(\lambda), M(\lambda)$, extra parameter $\lambda \in \mathbb{C}$
- ▷ Spectral curve: $0 = \det(L(\lambda) - \mu Id)$, does not depend on t

For this problem:

- $\mu^2 = Q_{h,k}(\lambda)$

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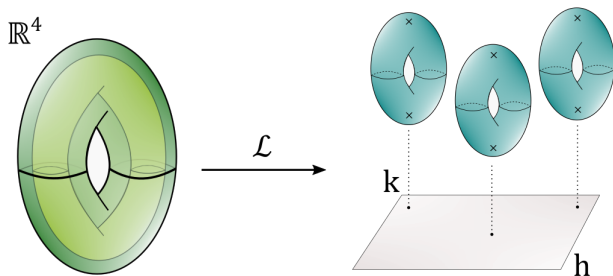
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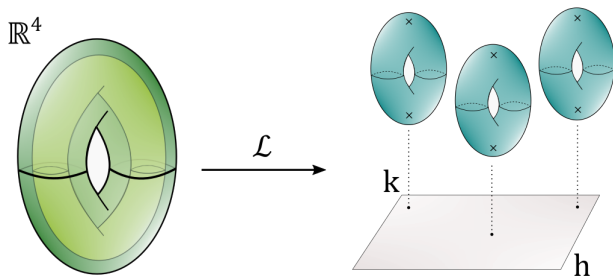
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Hamiltonian monodromy via spectral Lax pairs



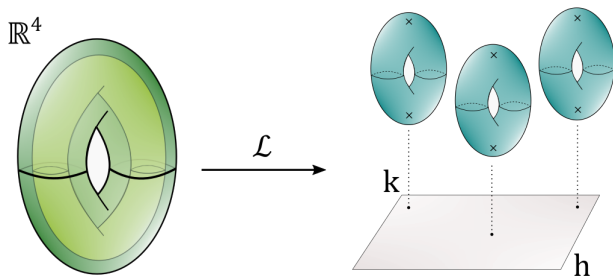
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Theorem [Mardesic, Sugny, G]

Consider the space (λ, μ) where $\mu^2 = Q_{h,k}(\lambda)$ and Q is a polynomial of degree 4 that has two non-real double roots for (h_0, k_0) . Let ξ be a 1-form in this space which can be written as $\xi = c_1 \frac{\lambda}{d\lambda} \mu + c_2 \frac{d\lambda}{\mu}$, $(c_1, c_2) \in \mathbb{C}^2$.

The variation of $\mathcal{I} = \int_{\delta} \xi$ when h and k vary along Γ is given by $\Delta_{\Gamma} \mathcal{I} = 2\pi i \operatorname{Res}(\xi, \infty)$. Moreover, this residue can be expressed as

$$\operatorname{Res}(\xi, \infty) = -\frac{c_1}{\sqrt{a_4}},$$

where a_4 is the leading coefficient of Q .

Corollary

Consider a completely integrable system in a neighborhood of an isolated critical value (h_0, k_0) with invariance under rotations associated to K and described by a spectral Lax pair with spectral curve $\mu^2 = Q_{h,k}(\lambda)$. Assume that the residue of η at infinity is $\frac{1}{i}$.

Under conditions given in the previous theorem, the monodromy matrix is

$$\mathcal{M}_\Gamma = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}$$

Spherical Pendulum

- ◇ Isolated critical value: $(h, k) = (1, 0)$
- ◇ Spectral Lax pair
- ◇ Spectral curve: $\mu^2 = \lambda^4 - 2k\lambda^3 + 2h\lambda^2 + 1$
- ◇ 1-form: $\eta = \frac{i\lambda}{\mu}d\lambda$, residue $\frac{1}{i}$ at infinity

$$\mathcal{M}_\Gamma = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}$$

non trivial HM \Rightarrow non global A-A coordinates

G. J. Gutierrez Guillen, P. Mardesic, D. Sugny, *Hamiltonian Monodromy via spectral Lax pairs.*

Thank you for the attention!